

Handedness inside the proton

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Questions to be addressed:

Is there nonzero transversity of quarks inside *unpolarized* hadrons?

How would one be able to find this out?

The transverse polarization of a *noncollinear* quark inside an unpolarized hadron in principle can have a preferred direction

This preferred direction signals an *intrinsic handedness*

Why? For instance, in the infinite momentum frame:

$$S_T^q \sim P_{\text{hadron}} \times p_{\text{quark}}$$

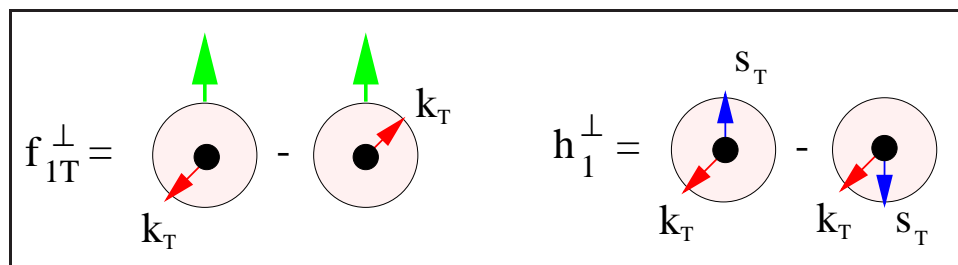
Obviously related to *orbital angular momentum*, but how exactly is still an open question

T-odd distribution functions

Such handedness appears to violate time reversal invariance
It is described by a **T-odd distribution function**, which was thought to be **absent** if the incoming hadron is treated as a **plane-wave state**

Recent work by Brodsky, Hwang, Schmidt, PLB 530 (2002) 99; Collins, PLB 536 (2002) 43; Ji, Yuan, hep-ph/0206057; Belitsky, Ji, Yuan, hep-ph/0208038 **shows otherwise**

Two leading twist (unsuppressed) T-odd distribution functions **with transverse momentum dependence** are possible



$f_{1T}^\perp \leftrightarrow$ **Sivers effect** (Sivers, PRD 41 (1990) 83; 43 (1991) 261)

h_1^\perp signals an **intrinsic handedness**

$$h_1^\perp \neq 0 \leftrightarrow \text{Prob}[q(\mathbf{k}_T, \mathbf{s}_T) \text{ in } P] \neq \text{Prob}[q(\mathbf{k}_T, -\mathbf{s}_T) \text{ in } P]$$

Its phenomenology was presented in D.B. & Mulders, PRD 57 (1998) 5780 and D.B., PRD 60 (1999) 014012 & NPB (PS) 79 (1999) 638

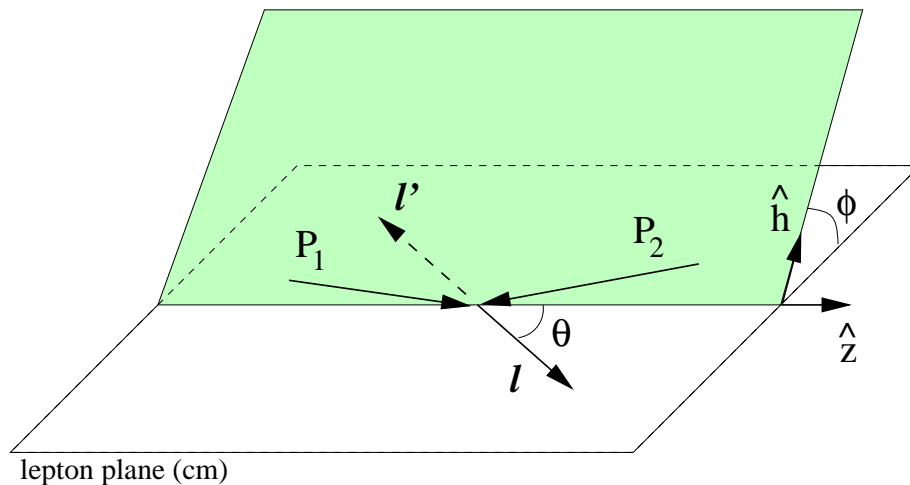
The unpolarized Drell-Yan process

There exists data compatible with nonzero h_1^\perp

Data from: NA10 Collab. ('86/'88); E615 Collab. at Fermilab ('89)
 $\pi^- N \rightarrow \mu^+ \mu^- X$, with $N = D, W$ and π -beams of 140-286 GeV

$$\frac{1}{\sigma} \frac{d\sigma}{d\Omega} \propto \left(1 + \lambda \cos^2 \theta + \mu \sin^2 \theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi \right)$$

Perturbative QCD prediction (NLO): $\lambda \approx 1$, $\mu \approx 0$, $\nu \approx 0$



Data: large ν !

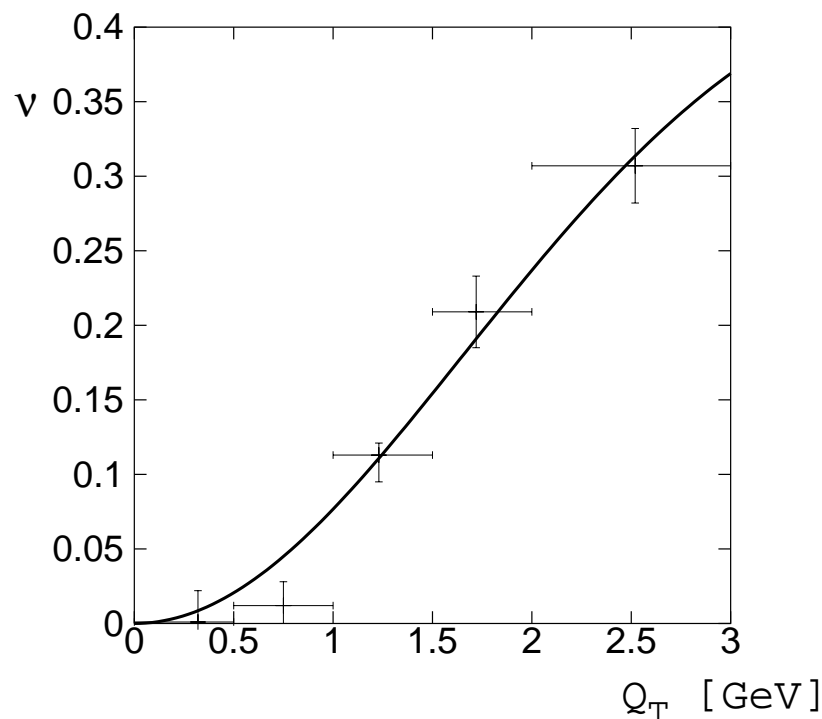
The unpolarized Drell-Yan process

The function h_1^\perp can provide an explanation for this large $\cos 2\phi$

Brandenburg, Nachtmann & Mirkes (ZPC 60 (1993) 697): large ν arises from a factorization breaking correlation between π and N

Observation: $\nu \propto h_1^\perp(\pi) h_1^\perp(N)$ [D.B., PRD 60 (1999) 014012]

Use the data to fit the function h_1^\perp



No factorization breaking and offers a natural explanation for $\mu \approx 0$

The polarized Drell-Yan process

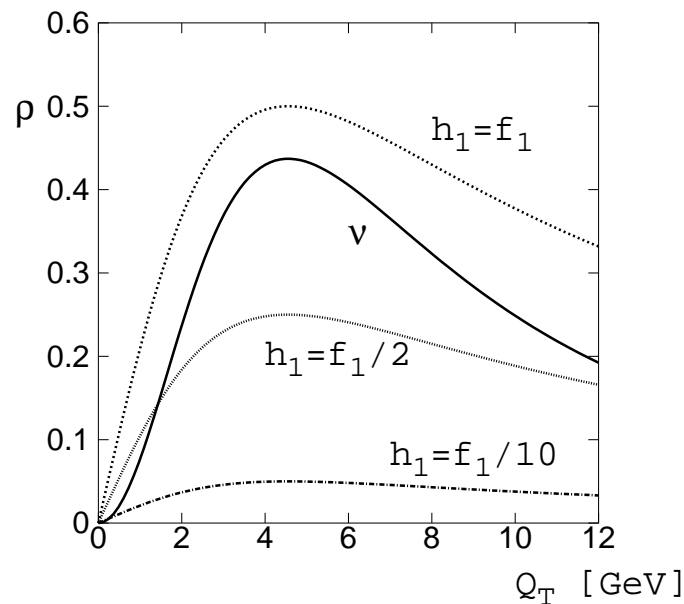
In the case of one **polarized hadron** (choosing $\mu = 0$ and $\lambda = 1$):

$$\frac{d\sigma}{d\Omega d\phi_S} \propto 1 + \cos^2 \theta + \sin^2 \theta \left[\frac{\nu}{2} \cos 2\phi - \rho |S_{1T}| \sin(\phi + \phi_S) \right] + \dots$$

Relation for the case of one flavor:

$$\rho = \frac{1}{2} \sqrt{\frac{\nu}{\nu_{\max}}} \frac{h_1}{f_1}$$

Yields as crude predictions



Different angular dependence compared to the **Sivers asymmetry**

$$(1 + \cos^2 \theta) |S_T| \sin(\phi - \phi_S) f_{1T}^\perp f_1$$

Further remarks

Possible future DY data

RHIC: $\langle \cos 2\phi \rangle$ in unpolarized $pp \rightarrow \mu^+ \mu^- X$ and the single spin asymmetry $\langle \sin(\phi + \phi_S) \rangle$, which is proportional to $h_1 h_1^\perp$

Fermilab: $\langle \cos 2\phi \rangle$ in $p\bar{p} \rightarrow \mu^+ \mu^- X$ probably yields larger results

Semi-inclusive DIS

The $\langle \cos 2\phi \rangle$ in SIDIS seems to be smaller than the $\langle \cos 2\phi \rangle$ in DY, but also smaller than the $\langle \cos \phi \rangle$ indicating that hard gluon radiation could be the explanation in SIDIS

In the present picture the $\langle \cos 2\phi \rangle$ in SIDIS would be $\propto h_1^\perp H_1^\perp$, hence this could be a sign that the magnitude of H_1^\perp is smaller than that of h_1^\perp

Testable by comparing to $\langle \cos 2\phi \rangle$ in e^+e^- annihilation (BELLE, BABAR)

Another test would be to look at $\langle \cos 2\phi \rangle$ for a jet instead of a hadron, $ep \rightarrow e' \text{jet } X$; h_1^\perp should then not contribute

SSA in hadron-hadron collisions

Single spin asymmetries in $p + p^\uparrow \rightarrow \pi + X$

These can arise from leading twist T-odd functions with transverse momentum dependence in three different ways:

- Distribution functions: $f_{1T}^\perp(x_1, \mathbf{p}_T) \otimes f_1(x_2) \otimes D_1(z)$
 $h_1^\perp(x_1, \mathbf{p}_T) \otimes h_1(x_2) \otimes D_1(z)$
- Fragmentation functions: $h_1(x_1) \otimes f_1(x_2) \otimes H_1^\perp(z, \mathbf{k}_T)$

Options 1 & 3 investigated by Anselmino, Boglione & Murgia (PLB 362 (95) 164 & PRD 60 (99) 054027)

The Collins effect, $H_1^\perp(z, \mathbf{k}_T)$, is expected to be present in quark fragmentation, but its **magnitude** is in principle **unrelated to** $h_1^\perp(x_1, \mathbf{p}_T)$

Options 1 & 2 also occur in $p + p^\uparrow \rightarrow \text{jet} + X$

Polarized Λ production in SIDIS

Polarized Λ production in semi-inclusive DIS

The intrinsic handedness can also lead to the following asymmetries:

- $\sin(\phi_{\Lambda}^e + \phi_{S_T^{\Lambda}}^e)$ asymmetry in $ep \rightarrow e'\Lambda^{\uparrow}X$
- $\sin(2\phi_{\Lambda}^e)$ asymmetry in $ep \rightarrow e'\vec{\Lambda}X$

These distinct angular dependences should be **absent** for **charged current exchange processes**, like $\nu p \rightarrow e\Lambda^{\uparrow}X$ or $\nu p \rightarrow e\vec{\Lambda}X$
D.B., Jakob & Mulders, NPB 564 (2000) 471

Distinguishable from other mechanisms via y or ϕ^e dependence

E.g. the **polarizing fragmentation functions** in SIDIS

Anselmino, D.B., D'Alesio & Murgia, PRD 65 (2002) 114014

Moreover, it should vanish after integration over Q_T , leaving only a $\sin(\phi_{S_T^{\Lambda}}^e)$ asymmetry (**twist-3**)

Conclusions

Leading twist T-odd distribution functions with transverse momentum dependence can offer viable explanations of certain azimuthal asymmetries

The intrinsic handedness function h_1^\perp can explain:

- the $\cos 2\phi$ asymmetry in unpolarized Drell-Yan
- single spin asymmetries in $p + p^\uparrow \rightarrow \pi + X$

Furthermore, it can generate

- $\sin(\phi + \phi_S)$ in $pp^\uparrow \rightarrow \mu^+ \mu^- X$
- $\cos 2\phi$ in $ep \rightarrow e' \pi X$, but not $ep \rightarrow e' \text{jet } X$
- $\sin(\phi_\Lambda^e + \phi_{S_T^\Lambda}^e)$ in $ep \rightarrow e' \Lambda^\uparrow X$ (NC, but not CC)
- $\sin(2\phi_\Lambda^e)$ in $ep \rightarrow e' \vec{\Lambda} X$ (NC, but not CC)

There are several ways of differentiating between mechanisms

Testable (in principle) using a host of existing (Fermilab, BELLE, ...) and future data (RHIC, COMPASS, HERMES, ...)